

# HP-KB<sub>3</sub>O<sub>5</sub> Highlights the Structural Diversity of Borates: Corner-Sharing BO<sub>3</sub>/BO<sub>4</sub> Groups in Combination with Edge-Sharing BO<sub>4</sub> Tetrahedra

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The new compound HP-KB<sub>3</sub>O<sub>5</sub> was synthesized under high-pressure/high-temperature conditions. It is the first compound exhibiting all three possible conjunctions simultaneously: corner-sharing BO<sub>3</sub> groups, corner-sharing BO<sub>4</sub>

units, as well as edge-sharing BO<sub>4</sub> tetrahedra. Calculations of harmonic vibrational frequencies and <sup>11</sup>B solid-state NMR measurements were performed for the first time with a compound exhibiting edge-sharing BO<sub>4</sub> tetrahedra.

## Introduction

The vast majority of syntheses in solid-state chemistry are performed by the variation of the thermodynamic parameters, composition, and temperature. The technically much-more demanding manipulation of the parameter pressure opens up tremendous possibilities for the synthesis of new materials. In the last decade, our high-pressure research focused on the substance class of borates, which led to a variety of new polymorphs and compounds with new compositions, not attainable under ambient-pressure conditions.<sup>[1]</sup> In contrast to ambient-pressure borates, in which trigonal BO<sub>3</sub> groups occur next to BO<sub>4</sub> tetrahedra, the new high-pressure borates predominantly exhibit boron fourfold coordinated by oxygen anions as a result of the pressure coordination rule.<sup>[2,3]</sup> Most striking was the discovery of edge-sharing BO<sub>4</sub> tetrahedra (B<sub>2</sub>O<sub>6</sub> units) in Dy<sub>4</sub>B<sub>6</sub>O<sub>15</sub>.<sup>[4]</sup>

Meanwhile, this motif was observed in three different structure types represented by the high-pressure compounds RE<sub>4</sub>B<sub>6</sub>O<sub>15</sub> (RE = Dy, Ho),<sup>[4,5]</sup> α-RE<sub>2</sub>B<sub>4</sub>O<sub>9</sub> (RE = Sm–Ho),<sup>[6–8]</sup> and β/HP-MB<sub>2</sub>O<sub>4</sub> (M = Fe, Ni, Co).<sup>[9–11]</sup> With the compounds M<sub>6</sub>B<sub>22</sub>O<sub>39</sub>·H<sub>2</sub>O (M = Fe, Co), intermediate states between corner- and edge-sharing BO<sub>4</sub> tetrahedra could also be isolated.<sup>[12]</sup> To our astonishment, the structural feature of edge-sharing BO<sub>4</sub> tetrahedra was recently found in the compound KZnB<sub>3</sub>O<sub>6</sub>, synthesized under ambient pressure conditions.<sup>[13,14]</sup> Thus, the structural motif of edge-sharing BO<sub>4</sub> tetrahedra is no longer a domain of high-pressure chemistry, but is still favored under these conditions, as there are eleven different compounds with three structure types prepared under high-pressure conditions and only one compound synthesized under ambient-pressure conditions. However, the synthesis of KZnB<sub>3</sub>O<sub>6</sub><sup>[13,14]</sup> leads to the question, in principle, whether it is possible to synthesize a borate exhibiting all the above-mentioned structural motifs (corner-sharing BO<sub>3</sub>/BO<sub>4</sub> groups and edge-sharing BO<sub>4</sub> tetrahedra) simultaneously. With HP-KB<sub>3</sub>O<sub>5</sub> (HP = high pressure), we synthesized a new borate, in which the two well-known fundamental structural motifs, alias BO<sub>3</sub> and BO<sub>4</sub> groups, and all noted conjunctions, namely corner-sharing of BO<sub>3</sub>/BO<sub>4</sub> groups and corner- as well as edge-sharing of BO<sub>4</sub> tetrahedra, are discovered in one crystal structure for the first time. This occurred in the scope of pressure-induced crystallization experiments on glass forming alkali borates, whilst investigating the ternary system K–B–O, where six different crystal structures are known: K<sub>5</sub>B<sub>19</sub>O<sub>31</sub>,<sup>[15]</sup> KB<sub>5</sub>O<sub>8</sub>,<sup>[16,17]</sup> α-KB<sub>5</sub>O<sub>8</sub>,<sup>[18]</sup> K<sub>2</sub>B<sub>4</sub>O<sub>7</sub>,<sup>[19]</sup> K<sub>3</sub>B<sub>3</sub>O<sub>6</sub>,<sup>[20,21]</sup> and KB<sub>3</sub>O<sub>5</sub>.<sup>[22,23]</sup> The latter (ambient-pressure modification) is built up from rigid triborate rings consisting of two corner-sharing BO<sub>3</sub> triangles and a corner-sharing BO<sub>4</sub> tetrahedron.

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## Results and Discussion

We succeeded in synthesizing the high-pressure modification HP-KB<sub>3</sub>O<sub>5</sub> by a high-pressure/high-temperature reaction at 3 GPa and 600 °C in a modified Walker module. The compound was obtained as a colorless crystalline solid. The crystal structure of HP-KB<sub>3</sub>O<sub>5</sub> was solved and refined by single-crystal X-ray crystallography. The crystal structure (Figure 1; view along [110]) consists of ribbons of BO<sub>4</sub> tetrahedra, which run along the *c* axis.

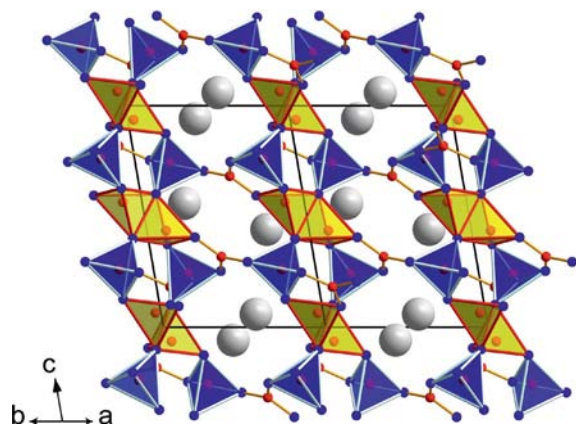


Figure 1. Crystal structure of HP-KB<sub>3</sub>O<sub>5</sub> with edge-sharing BO<sub>4</sub> units; view along [110]. Large light spheres: K<sup>+</sup>, corners of polyhedra: O<sup>2-</sup>, small grey spheres: B<sup>3+</sup>.

Within the ribbons, the BO<sub>4</sub> tetrahedra are connected through common corners and edges. The ribbons are interconnected by trigonal-planar BO<sub>3</sub> groups to form channels, in which the tenfold coordinated potassium ions are situated. The recently published KZnB<sub>3</sub>O<sub>6</sub><sup>[13,14]</sup> shows remarkable similarities to our compound and is built up from [B<sub>6</sub>O<sub>12</sub>]<sup>6-</sup> groups (two rings of two BO<sub>3</sub> and one BO<sub>4</sub> unit, linked by edge-sharing BO<sub>4</sub> tetrahedra), which are connected by ZnO<sub>4</sub> tetrahedra to form a three-dimensional network. Figure 2 depicts a comparison of the crystal structures of HP-KB<sub>3</sub>O<sub>5</sub> and KZnB<sub>3</sub>O<sub>6</sub><sup>[13,14]</sup>. In HP-KB<sub>3</sub>O<sub>5</sub> (Figure 2, left), the edge-sharing BO<sub>4</sub> tetrahedra (light polyhedra) connect to two BO<sub>3</sub> and four BO<sub>4</sub> groups through common corners. Two of the BO<sub>4</sub> groups are situated at the common edge of the B<sub>2</sub>O<sub>6</sub> unit. It is for the first time that oxygen ions, involved in the edge-sharing of BO<sub>4</sub> tetrahedra, are observed to bridge three BO<sub>4</sub> tetrahedra (O<sup>[3]</sup>). In KZnB<sub>3</sub>O<sub>6</sub><sup>[13,14]</sup> (Figure 2, right), there are also edge-sharing BO<sub>4</sub> tetrahedra and BO<sub>3</sub> groups, but no corner-sharing BO<sub>4</sub> groups. Here, the edge-sharing BO<sub>4</sub> tetrahedra (light polyhedra) connect to four BO<sub>3</sub> units; two ZnO<sub>4</sub> groups are situated at the common edge of the B<sub>2</sub>O<sub>6</sub> unit (highlighted with arrows).

The mean B–O distance in the BO<sub>4</sub> tetrahedra and the BO<sub>3</sub> groups of HP-KB<sub>3</sub>O<sub>5</sub> is 147.7 and 137.3 pm, respectively, which agrees with the average B–O distance of the BO<sub>4</sub> and BO<sub>3</sub> units in borates (BO<sub>4</sub>: 147.6 pm, BO<sub>3</sub>: 137.0 pm).<sup>[24–26]</sup> Figure 3 (left) shows that the B–O dis-

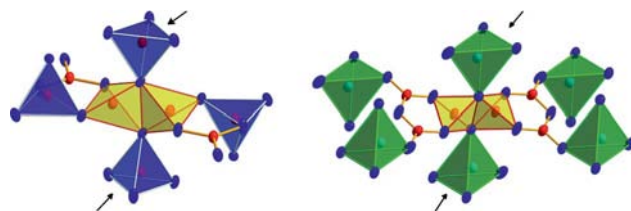


Figure 2. Comparison of the central structural motifs of HP-KB<sub>3</sub>O<sub>5</sub> (left) and KZnB<sub>3</sub>O<sub>6</sub><sup>[13,14]</sup> (right) (90% probability ellipsoids). The arrows point to the BO<sub>4</sub> (HP-KB<sub>3</sub>O<sub>5</sub>) and ZnO<sub>4</sub> tetrahedra (KZnB<sub>3</sub>O<sub>6</sub>), which are located at the common edge of the edge-sharing BO<sub>4</sub> groups.

tances inside the B<sub>2</sub>O<sub>2</sub> rings of the B<sub>2</sub>O<sub>6</sub> unit are larger [154.8(1) and 151.4(1) pm] than those outside the rings [141.2(1) and 144.6(1) pm]. The B···B distance inside the B<sub>2</sub>O<sub>6</sub> unit of HP-KB<sub>3</sub>O<sub>5</sub> is 221.5(1) pm, which is significantly larger than the values found in other high-pressure borates that exhibit edge-sharing BO<sub>4</sub> tetrahedra, e.g. Dy<sub>4</sub>B<sub>6</sub>O<sub>15</sub> [207.2(8) pm],<sup>[4]</sup> α-Sm<sub>2</sub>B<sub>4</sub>O<sub>9</sub> [207.1(9) pm],<sup>[8]</sup> β-FeB<sub>2</sub>O<sub>4</sub> [208.3(5) pm],<sup>[9]</sup> and KZnB<sub>3</sub>O<sub>6</sub> [207.9(4) pm].<sup>[13,14]</sup> The reason for this increased value lies in the different coordination spheres of the O ions at the common edge of the B<sub>2</sub>O<sub>6</sub> unit (O1). As depicted in Figure 3 (right), the O1 anions in β-FeB<sub>2</sub>O<sub>4</sub><sup>[9]</sup> are coordinated tetrahedrally in the form of two covalently bound boron atoms and two coordinating metal ions (2+2 coordination). In HP-KB<sub>3</sub>O<sub>5</sub> (Figure 3, left), the O1 anions are bound covalently to three boron atoms (O<sup>[3]</sup>), which lead to enhanced distances inside the B<sub>2</sub>O<sub>6</sub> ring and therewith to a remarkably enlarged B···B distance.

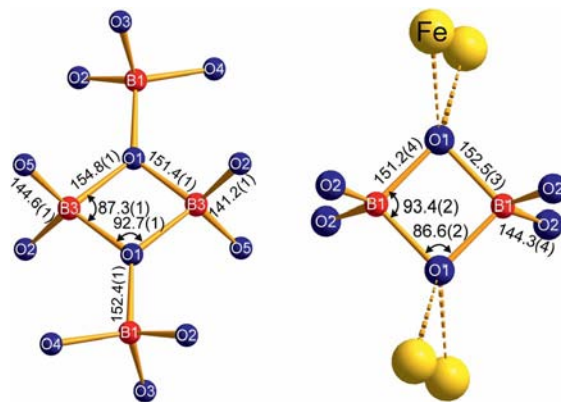


Figure 3. Comparison of selected distances [pm] and angles [°] in the B<sub>2</sub>O<sub>6</sub> group of HP-KB<sub>3</sub>O<sub>5</sub> (left) and β-FeB<sub>2</sub>O<sub>4</sub> (right).

The potassium ions are coordinated by ten oxygen ions with distances between 264–341 pm (av. 300 pm). These values correspond to those in KV<sub>3</sub>O<sub>8</sub> (272–331 pm; av. 301 pm)<sup>[27]</sup> and K<sub>3</sub>V<sub>5</sub>O<sub>14</sub> (273–335 pm; av. 300 pm).<sup>[28]</sup>

Bond-valence sums were calculated for all atoms by using the CHARDI (charge distribution in solids, ΣQ)<sup>[29]</sup> and bond length/bond strength concepts (ΣV).<sup>[30,31]</sup> The formal charges of the ions were confirmed by the calculations of both concepts.

Taking into account the synthetic conditions (3 GPa, 600 °C), we investigated the assumed metastable character of the high-pressure phase HP-KB<sub>3</sub>O<sub>5</sub>. Temperature-programmed X-ray powder diffraction experiments reveal that HP-KB<sub>3</sub>O<sub>5</sub> is stable up to 550 °C under ambient-pressure conditions (Figure S1). Higher temperatures lead inter alia to the normal pressure phase KB<sub>3</sub>O<sub>5</sub>.<sup>[22]</sup> DTA measurements confirm the temperature stability of HP-KB<sub>3</sub>O<sub>5</sub> (Figure S2).

Vibrational spectroscopic measurements and calculations of harmonic vibrational frequencies at the  $\Gamma$  point yield modes involving mainly the corner-sharing BO<sub>3</sub> groups above 1320 cm<sup>-1</sup>, stretching vibrations of the corner- and edge-sharing BO<sub>4</sub> tetrahedra between 950–1215 cm<sup>-1</sup>, bending and complex vibrations of both BO<sub>3</sub> and BO<sub>4</sub> units between 200–905 cm<sup>-1</sup>, and lattice vibrations involving potassium below 185 cm<sup>-1</sup> (Figure 4).

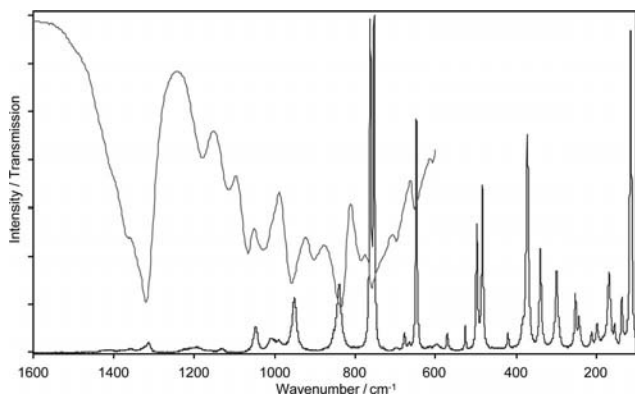


Figure 4. FTIR-ATR (top) and confocal Raman (bottom) spectra of HP-KB<sub>3</sub>O<sub>5</sub>. For a better comparability, the ATR-units were transformed into transmission.

The calculations enabled for the first time a precise assignment of the previously discussed vibrational modes of these B<sub>2</sub>O<sub>6</sub> groups.<sup>[32,12]</sup> For example, bands at 1213, 1205, 1161, and 1013 cm<sup>-1</sup> for the Raman-active modes and at 1105, 1070, and 1001 cm<sup>-1</sup> for the IR-active modes can now be assigned to modes involving strong movements of the two common oxygen ions within the two edge-sharing tetrahedra. The Raman-active mode at 1213 cm<sup>-1</sup> exhibits the strongest vibration within the B<sub>2</sub>O<sub>6</sub> unit.

The Raman-active modes of B<sub>2</sub>O<sub>6</sub> groups were investigated previously on HP-NiB<sub>2</sub>O<sub>4</sub>,<sup>[9]</sup>  $\alpha$ -Gd<sub>2</sub>B<sub>4</sub>O<sub>9</sub>,<sup>[6]</sup> and Dy<sub>4</sub>B<sub>6</sub>O<sub>15</sub>.<sup>[4]</sup> The spectra of these compounds reveal several bands in the range 1200–1450 cm<sup>-1</sup>, which normally correspond to BO<sub>3</sub> or OB<sub>3</sub> groups (the latter is not part of a B<sub>2</sub>O<sub>6</sub> group). As none of the three compounds with edge-sharing BO<sub>4</sub> tetrahedra contain three-coordinate boron atoms or three-coordinate oxygen atoms, these peaks are assumed to likely be Raman-active modes of the B<sub>2</sub>O<sub>6</sub> unit. A comparison of the different spectra indicates two bands that probably could be assigned to modes of edge-sharing BO<sub>4</sub> tetrahedra (1262 and 1444 cm<sup>-1</sup> in HP-NiB<sub>2</sub>O<sub>4</sub>, 1253 and 1431 cm<sup>-1</sup> in  $\alpha$ -Gd<sub>2</sub>B<sub>4</sub>O<sub>9</sub>, and 1271 and 1435 cm<sup>-1</sup> in

Dy<sub>4</sub>B<sub>6</sub>O<sub>15</sub>). Even for the recently published KZnB<sub>3</sub>O<sub>6</sub>,<sup>[13,14]</sup> Jin et al. found an absorption band in the range 1400 cm<sup>-1</sup>, but missed the weak B<sub>2</sub>O<sub>6</sub> vibration at 1210 cm<sup>-1</sup>. Further calculations are necessary to reveal whether the B<sub>2</sub>O<sub>6</sub> units of the above-mentioned phases show bands in the region of 1400 cm<sup>-1</sup>. The here discussed HP-KB<sub>3</sub>O<sub>5</sub> exhibits no band around 1400 cm<sup>-1</sup>, which is caused by the vibrations of the B<sub>2</sub>O<sub>6</sub> unit. This difference in the spectra from those of other compounds comprising edge-sharing BO<sub>4</sub> tetrahedra might be due to the increased boron coordination of O1. Vibrational spectroscopy of HP-KB<sub>3</sub>O<sub>5</sub> confirms the absence of water in the structure.

Until 2009, only paramagnetic compounds with edge-sharing BO<sub>4</sub> tetrahedra could be obtained, on which solid-state <sup>11</sup>B NMR spectroscopic investigations are impossible. Such <sup>11</sup>B solid-state NMR investigations would be of paramount interest to answer the question whether the chemical shift values and quadrupolar coupling constants of <sup>11</sup>B are sensitive enough to unequivocally identify the structural motif of the edge-sharing BO<sub>4</sub> tetrahedra. This would significantly simplify the investigation of technically important borate glasses. Up to now, edge-sharing BO<sub>4</sub> tetrahedra are unknown in borate glass chemistry.<sup>[33,34]</sup> With the two diamagnetic compounds HP-KB<sub>3</sub>O<sub>5</sub> and KZnB<sub>3</sub>O<sub>6</sub>,<sup>[13,14]</sup> <sup>11</sup>B solid-state NMR measurements now become possible for the first time.

Figure 5 shows a <sup>11</sup>B MQMAS spectrum for HP-KB<sub>3</sub>O<sub>5</sub>. The MQMAS technique allows the separation of the individual resonances of the typical broad 1D <sup>11</sup>B MAS spectra. While the F2 dimension retains the full interaction including the chemical shift and second-order quadrupole interaction, for the F1 dimension, only the isotropic chemical and quadrupolar induced shifts are relevant. For HP-KB<sub>3</sub>O<sub>5</sub>, both inequivalent BO<sub>4</sub> as well as the BO<sub>3</sub> units can be distinguished (Table 1). Furthermore, an evaluation of slices at  $\delta(F1)$  values of 0.5, 2.5, and 22 ppm gives access to the isotropic chemical shift  $\delta_{iso}$ , the quadrupole coupling constant  $C_Q$ , and the asymmetry parameter  $\eta_Q$  for each individual boron site (Figure 5 and Table 1). These data were then used to determine the intensity ratios by deconvoluting an 1D <sup>11</sup>B MAS spectrum recorded with selective excitation (Figure 5). In a similar procedure, the data for KZnB<sub>3</sub>O<sub>6</sub><sup>[13,14]</sup> depicted in Table 1 were also collected. Although data analysis of KZnB<sub>3</sub>O<sub>6</sub> was hampered by a substantial impurity of boronic acid, values for  $\delta_{iso}$ ,  $C_Q$ , and  $\eta_Q$  for the trigonal planar and tetrahedral boron units could be derived.

On the basis of investigations of minerals, some empirical values are reported for both BO<sub>3</sub> and BO<sub>4</sub> units.<sup>[35–38]</sup> While BO<sub>3</sub> units typically are observed between 14 and 23 ppm ( $\delta_{iso}$ ) and exhibit large  $C_Q$  values (2.4 to 2.9 MHz), tetrahedrally coordinated boron atoms usually range from –1 to 3.5 ppm and 0.2 to 0.5 MHz for  $\delta_{iso}$  and  $C_Q$ , respectively. The asymmetry parameter  $\eta$  largely depends on local distortions and connectivities and can vary between 0 and 0.8 for both units. The observed values (Table 1) for HP-KB<sub>3</sub>O<sub>5</sub> and KZnB<sub>3</sub>O<sub>6</sub> fit quite well into these regions. Since for KZnB<sub>3</sub>O<sub>6</sub><sup>[13,14]</sup> only edge-sharing BO<sub>4</sub> units exist, the iso-



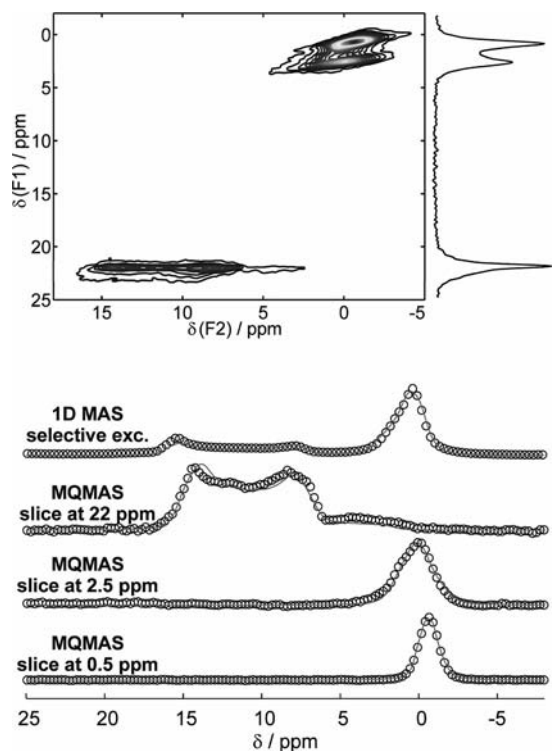


Figure 5.  $^{11}\text{B}$  MAS NMR spectra for  $\text{HP-KB}_3\text{O}_5$ . 2D MQMAS spectrum (top) and 1D MAS spectrum with selective excitation as well as three horizontal slices of the MQMAS spectrum (bottom). The experimental data and the simulations are depicted with open circles and solid lines, respectively. The relevant refinement parameters are given in Table 1.

Table 1. Results of the refinement of the  $^{11}\text{B}$  MQMAS and selective 1D spectra for  $\text{HP-KB}_3\text{O}_5$  and  $\text{KZnB}_3\text{O}_6$ . Details of the fitting procedure are given in the text.

	Unit	$\delta_{\text{iso}} / \text{ppm}$	$C_Q / \text{MHz}$	$\eta$	$I$
HP-KB <sub>3</sub> O <sub>5</sub>	BO <sub>4</sub>	0.31(2)	0.95(5)	0.55(4)	1.08(6)
	BO <sub>4</sub>	1.75(2)	1.2(1)	0.74(5)	0.85(6)
	BO <sub>3</sub> (T <sub>2</sub> )	16.8(2)	2.5(1)	0.17(3)	0.98(6)
KZnB <sub>3</sub> O <sub>6</sub>	BO <sub>4</sub>	2.4(1)	0.8(1)	0.4(1)	–
	BO <sub>3</sub> (T <sub>2</sub> )	18.1(3)	2.4(1)	0.3(2)	–

tropic shift of 2.4 ppm can be assigned to those units. Probably, the resonance at  $\delta = 1.75$  ppm might accordingly be assigned to the edge-sharing BO<sub>4</sub> tetrahedra for  $\text{HP-KB}_3\text{O}_5$ . Nevertheless, because there are only slight variations for the second BO<sub>4</sub> unit, this assignment is not unequivocal. It is at least questionable whether it will be possible to detect edge-sharing tetrahedral units on the basis of  $^{11}\text{B}$  chemical shift data. In a forthcoming study, we will combine NMR spectroscopic analysis with quantum chemical calculations of the chemical shift and the electric field gradient tensor to get a more detailed picture of the relevant interactions.

## Conclusions

Systemic investigations in the field of high-pressure borates led to the synthesis of  $\text{HP-KB}_3\text{O}_5$ , the first borate exhibiting all three possible conjunctions: corner-sharing BO<sub>3</sub> groups, corner-sharing BO<sub>4</sub> units, as well as edge-sharing BO<sub>4</sub> tetrahedra. Vibrations of the B<sub>2</sub>O<sub>6</sub> unit could be assigned in a measured spectrum through calculations of the harmonic vibrational frequencies at the  $\Gamma$  point. With the two diamagnetic compounds  $\text{HP-KB}_3\text{O}_5$  and  $\text{KZnB}_3\text{O}_6$ ,  $^{11}\text{B}$  solid-state NMR-measurements were performed for the first time for compounds exhibiting edge-sharing BO<sub>4</sub> tetrahedra.

## Experimental Section

**Synthesis:** The title compound was prepared by grinding a stoichiometric mixture of K<sub>2</sub>CO<sub>3</sub> (ChemPur GmbH, Karlsruhe, Germany, 99.9%) and H<sub>3</sub>BO<sub>3</sub> (Merck KGaA, Darmstadt, Germany, 99.5%) and heating it at 900 °C for 12 h to expel H<sub>2</sub>O and CO<sub>2</sub>. To obtain  $\text{HP-KB}_3\text{O}_5$ , the amorphous reaction product was compressed and heated in a multianvil assembly. Therefore, the sample was re-ground and filled into a boron nitride crucible (Henze BNP GmbH, HeBoSint® S10, Kempten, Germany). The crucible was placed into the pressure medium: a MgO octahedron (Ceramic Substrates & Components, Isle of Wight, UK) with an edge length of 18 mm. The octahedron was compressed via eight tungsten carbide cubes (TSM-10, Ceratizit, Reutte, Austria) with a truncation edge length of 11 mm (18/11 assembly). Pressure was applied by a Walker-type multianvil device and a 1000 t press (both devices from the company Voggenreiter, Mainleus, Germany). A detailed description of the assembly and its preparation can be found in ref.<sup>[39–43]</sup> For the synthesis of  $\text{HP-KB}_3\text{O}_5$ , the assembly was compressed to 3 GPa over 1.25 h and kept at this pressure for the heating period. The sample was heated to 600 °C over 10 min, kept there for 10 min, and cooled down to 400 °C over 40 min. Afterwards, the sample was naturally cooled down to room temperature by switching off the heating, followed by a decompression period of 3.75 h. The recovered pressure medium was broken apart, and the surrounding boron nitride crucible removed from the sample. The compound  $\text{HP-KB}_3\text{O}_5$  was obtained as an air-resistant, colorless crystalline solid.

**X-ray Structure Determination:** Crystal structure analysis of  $\text{HP-KB}_3\text{O}_5$ :  $M_w = 151.53 \text{ g mol}^{-1}$ , colorless block,  $0.12 \times 0.11 \times 0.06 \text{ mm}$ , monoclinic, space group:  $C2/c$ ,  $a = 9.608(2)$ ,  $b = 8.770(2)$ ,  $c = 9.099(2) \text{ \AA}$ ,  $\beta = 104.4(1)^\circ$ ,  $V = 742.8(3) \text{ \AA}^3$ ,  $Z = 8$ ,  $\rho_{\text{calcd.}} = 2.710 \text{ g cm}^{-3}$ , Nonius Kappa CCD, Mo-K $\alpha$  radiation ( $\lambda = 0.71073 \text{ \AA}$ ); graphite monochromator,  $F(000) = 592$ ,  $\mu = 1.33 \text{ mm}^{-1}$ ,  $T = 293(2) \text{ K}$ , multiscan, 1647 measured reflections in the range  $3.2 < 2\theta < 35.0^\circ$ , 1502 unique reflections with  $I < 2\sigma(I)$ ,  $R_{\text{int}} = 0.0237$ ; the crystal structure was solved by direct methods (SHELXS-97) and anisotropically refined by a least-squares procedure against  $F^2$  with all data, 82 refined parameters,  $R_1 = 0.0266$  and  $wR_2 = 0.0635$  for  $F_o^2 > 2(F_o^2)$ ;  $R_1 = 0.0312$  and  $wR_2 = 0.0650$  for all  $F_o^2$ , residual electron density between  $-0.61$  and  $0.44 \text{ e \AA}^{-3}$ , GOF = 1.107. Further details on the crystal structure investigations may be obtained from the Fachinformationszentrum Karlsruhe (FIZ, 76344 Eggenstein-Leopoldshafen, Germany; Fax: +49-7247-808-666; E-mail: crysdata@fiz-karlsruhe.de), on quoting the depository number CSD-423027.

**Temperature-Programmed X-ray Powder Diffraction:** Temperature-programmed X-ray powder diffraction experiments were performed

on a STOE Stadi P powder diffractometer (Mo- $K_{\alpha}$ ) with a computer controlled STOE furnace. The sample was enclosed in a silica-glass capillary and heated from room temperature to 500 °C in steps of 100 °C and from 500 to 1100 °C in steps of 50 °C. After reaching 1100 °C, the sample was cooled down to 500 °C in steps of 50 °C and further on to room temperature in steps of 100 °C. At each step, a powder diffraction pattern was recorded.

**DTA/TG:** The simultaneous differential thermal analysis and thermogravimetric (DTA/TG) measurements were performed with a Setaram TG-92 simultaneous thermal analyzer equipped with a protected DTA/TG rod (P-type thermocouple, maximum operating temperature = 1600 °C). Temperature calibration was achieved in reference to melting points of metal standards. The measurements were conducted in a helium atmosphere.

**Calculations of the Vibrational Frequencies:** The vibrational frequencies were computed from first principles, by using the program CRYSTAL 06,<sup>[44]</sup> Gaussian basis sets, and the B3LYP Hybrid HF-DFT functional.

**IR and Raman Spectroscopy:** The confocal Raman spectra of single crystals in the range 100–1500 cm<sup>-1</sup> were obtained with a HORIBA JOBIN YVON LabRam-HR 800 spectrometer. The sample was excited by the 532 nm emission lines of a 30 mW Nd:YAG laser. The size and power of the laser spot on the surface were approximately 1 µm and 5 mW. The spectral resolution, determined by measuring the Rayleigh line, was 1.4 cm<sup>-1</sup>. Raman bands were fitted by the built-in spectrometer software LabSpec 5 to convoluted Gauss-Lorentz functions. The accuracy of Raman line shifts was in the order of 0.5 cm<sup>-1</sup>. The FTIR-ATR (Attenuated Total Reflection) spectra of single crystals in the range 600–4000 cm<sup>-1</sup> were recorded with a Bruker Vertex 70 FTIR spectrometer (spectral resolution 4 cm<sup>-1</sup>), attached to a Hyperion 3000 microscope. A germanium ATR-crystal with a tip diameter of 100 µm was pressed on the surface of the borate crystal with a power of 5 N, which crushed it into pieces that were µm-sized. Spectra were corrected for atmospheric influences, and an enhanced ATR correction, with the OPUS 6.5 software, was performed. Background correction and peak fitting followed via polynomial and folded Gaussian-Lorentzian functions.

**<sup>11</sup>B Solid-State NMR Spectroscopy:** The <sup>11</sup>B solid-state NMR spectra were acquired under magic angle spinning conditions with a Bruker Avance 500 spectrometer operating at a resonance frequency of 160.461 MHz. The peak positions are given with respect to BF<sub>3</sub>(OEt<sub>2</sub>) at 0 ppm. The 1D spectra were recorded with a 2.5 mm triple resonance MAS probe (Bruker) by applying three soft back-to-back 90 ° pulses ( $\nu_{\text{nut}} = 10$  kHz) to ensure selective excitation, while simultaneously suppressing unwanted resonances from the probe.<sup>[45]</sup> Spinning speed and recycle delay were set to 25 kHz and 40 s, respectively. To collect the MQMAS spectra, we worked with a 3.2 triple resonance MAS probe (Bruker) and a spinning speed of 20 kHz. A three-pulse sequence with two leading hard pulses ( $\nu_{\text{nut}} = 180$  kHz) and a selective 90 ° pulse ( $\nu_{\text{nut}} = 10$  kHz) was used.<sup>[46]</sup> The pulse lengths were optimized to 3.4, 1.3, and 25 µs. The 30-fold phase cycle ensured that only the desired coherence pathway  $0 \rightarrow \pm 3 \rightarrow 0 \rightarrow -1$  was selected. To obtain pure absorptive spectra, we invoked the States scheme. In total, 178 fids with 60 repetitions were recorded by using a recycle delay of 6 s. The simulations of the 1D spectra and the slices of the MQMAS spectra were carried out with the program package SIMPSON.<sup>[47]</sup>

**Supporting Information** (see footnote on the first page of this article): Temperature-programmed powder pattern as well as DTA curve of HP-KB<sub>3</sub>O<sub>5</sub>.

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- [1] H. Huppertz, *Chem. Commun.* **2011**, 47, 131–140.
- [2] A. Neuhaus, *Chimia* **1964**, 18, 93–103.
- [3] C. T. Prewitt, R. T. Downs, *Rev. Mineral.* **1998**, 37, 283–317.
- [4] H. Huppertz, B. von der Eltz, *J. Am. Chem. Soc.* **2002**, 124, 9376–9377.
- [5] H. Huppertz, *Z. Naturforsch. B* **2003**, 58, 278.
- [6] H. Emme, H. Huppertz, *Chem. Eur. J.* **2003**, 9, 3623–3633.
- [7] H. Emme, H. Huppertz, *Z. Anorg. Allg. Chem.* **2002**, 628, 2165.
- [8] H. Emme, H. Huppertz, *Acta Crystallogr., Sect. C* **2005**, 61, i29.
- [9] S. C. Neumair, R. Glaum, H. Huppertz, *Z. Naturforsch. B* **2009**, 64, 883–890.
- [10] J. S. Knyrim, F. Roeßner, S. Jakob, D. Johrendt, I. Kinski, R. Glaum, H. Huppertz, *Angew. Chem.* **2007**, 119, 9256–9259; *Angew. Chem. Int. Ed.* **2007**, 46, 9097–9100.
- [11] S. C. Neumair, R. Kaindl, H. Huppertz, *Z. Naturforsch. B* **2010**, 65, 1311–1317.
- [12] S. C. Neumair, J. S. Knyrim, O. Oeckler, R. Glaum, R. Kaindl, R. Stalder, H. Huppertz, *Chem. Eur. J.* **2010**, 16, 13659–13670.
- [13] S. Jin, G. Cai, W. Wang, M. He, S. Wang, X. Chen, *Angew. Chem.* **2010**, 122, 5087–5090; *Angew. Chem. Int. Ed.* **2010**, 49, 4967–4970.
- [14] Y. Wu, J.-Y. Yao, J.-X. Zhang, P.-Z. Fu, Y.-C. Wu, *Acta Crystallogr., Sect. E* **2010**, 66, i45.
- [15] J. Krogh-Moe, *Acta Crystallogr., Sect. B* **1974**, 30, 1827–1832.
- [16] J. Krogh-Moe, *Acta Crystallogr.* **1965**, 18, 1088–1089.
- [17] J. Krogh-Moe, *Ark. Kemi.* **1959**, 14, 439–449.
- [18] J. Krogh-Moe, *Acta Crystallogr., Sect. B* **1972**, 28, 168–172.
- [19] J. Krogh-Moe, *Acta Crystallogr., Sect. B* **1972**, 28, 3089–3093.
- [20] W. H. Zachariasen, *J. Chem. Phys.* **1937**, 5, 919–922.
- [21] W. Schneider, G. B. Carpenter, *Acta Crystallogr., Sect. B* **1970**, 26, 1189–1191.
- [22] R. S. Bubnova, V. S. Fundamenskii, S. K. Filatov, I. G. Polyakova, *Dokl. Akad. Nauk* **2004**, 398, 643–647.
- [23] J. Krogh-Moe, *Acta Crystallogr.* **1961**, 14, 68.
- [24] E. Zobel, *Z. Kristallogr.* **1990**, 191, 45–57.
- [25] F. C. Hawthorne, P. C. Burns, J. D. Grice in “The Crystal Chemistry of Boron”, *Boron: Mineralogy, Petrology and Geochemistry*, Vol. 33, 2nd ed. (Eds.: E. S. Grew, L. M. Anovitz), Mineralogical Society of America, Washington, **1996**, pp. 41.
- [26] E. Zobel, *Z. Kristallogr.* **1982**, 160, 81–92.
- [27] Y. Oka, T. Yao, N. Yamamoto, *Mater. Res. Bull.* **1997**, 32, 1201–1209.
- [28] A. M. Byström, H. T. Evans Jr., *Acta Chem. Scand.* **1959**, 13, 377–378.
- [29] R. Hoppe, S. Voigt, H. Glaum, J. Kissel, H. P. Müller, K. J. Bernet, *J. Less-Common Met.* **1989**, 156, 105–122.
- [30] I. D. Brown, D. Altermatt, *Acta Crystallogr., Sect. B* **1985**, 41, 244–247.
- [31] N. E. Brese, M. O’Keeffe, *Acta Crystallogr., Sect. B* **1991**, 47, 192–197.
- [32] H. Huppertz, G. Heymann, *Solid State Sci.* **2003**, 5, 281–289.
- [33] A. C. Wright, *Phys. Chem. Glasses* **2010**, 51, 1–39.
- [34] A. C. Wright, G. Dalba, F. Rocca, N. M. Vedishcheva, *Phys. Chem. Glasses* **2010**, 51, 233–265.
- [35] S. Kroeker, J. F. Stebbins, *Inorg. Chem.* **2001**, 40, 6239–6246.
- [36] D. Müller, A. R. Grimmer, U. Timper, G. Heller, M. Shakibaie-Moghadam, *Z. Anorg. Allg. Chem.* **1993**, 619, 1262–1268.
- [37] M. R. Hansen, T. Vosegaard, H. J. Jakobsen, J. Skibs, *J. Phys. Chem. A* **2004**, 108, 586–594.
- [38] M. R. Hansen, G. K. H. Madsen, H. J. Jakobsen, J. Skibs, *J. Phys. Chem. A* **2005**, 109, 1989–1997.

- [39] N. Kawai, S. Endo, *Rev. Sci. Instrum.* **1970**, *41*, 1178–1181.
- [40] D. Walker, M. A. Carpenter, C. M. Hitch, *Am. Mineral.* **1990**, *75*, 1020–1028.
- [41] D. Walker, *Am. Mineral.* **1991**, *76*, 1092–1100.
- [42] D. C. Rubie, *Phase Transitions* **1999**, *68*, 431–451.
- [43] H. Huppertz, *Z. Kristallogr.* **2004**, *219*, 330–338.
- [44] R. Dovesi, V. R. Saunders, C. Roetti, R. Orlando, C. M. Zicovich-Wilson, F. Pascale, B. Civalleri, K. Doll, N. M. Harrison, I. J. Bush, Ph. D'Arco, M. Llunell, *CRYSTAL06 - User's Manual*, University of Torino, Torino, **2006**.
- [45] S. Zhang, X. L. Wu, M. Mehring, *Chem. Phys. Lett.* **1990**, *173*, 481–484.
- [46] J. P. Amoureux, C. Fernandez, S. Steuernagel, *J. Magn. Reson., Ser. A* **1996**, *123*, 116–118.
- [47] M. Bak, J. T. Rasmussen, N. C. Nielsen, *J. Magn. Reson.* **2000**, *147*, 296–330.

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